

Helicity of the W Boson in Lepton+Jets $t\bar{t}$ Events



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- Introduction
- The new approach for measuring top quark properties
- Monte Carlo tests with the new approach
- F_0 measurement using Run I DØ data
- Systematic uncertainties
- Conclusions

Event Topology and Selection Criteria

- **DØ Statistics Run I (125 pb⁻¹)**
 - In proton antiprotons collisions @ $\sqrt{s}=1.8\text{TeV}$ top quarks are primarily **produced in pairs**
 - **90% qqbar $\rightarrow t\bar{t}$, 10% gg- $t\bar{t}$**
 - Each top quark decays weakly: **BR($t \rightarrow Wb$) @ 100%**
 - There are 3 main experimental $t\bar{t}$ signatures depending on the decay of the W bosons:
 - Dilepton BR($e\bar{e}+\mu\bar{\mu}+e\mu$) ~ 5%
 - All-Hadronic BR(quarks) ~ 44%
 - **Lepton +Jets BR($e+\text{jets}, \mu+\text{quarks}$) ~ 30%**

Lepton+Jets

- **Signal:** 1 high- P_T lepton, 4 jets, large missing- E_T
- **Background:** W with associated production of jets

Standard Selection:

- **Lepton:** $E_T > 20\text{GeV}$, $|\eta^e| < 2$, $|\eta^\mu| < 1.7$
- **Jets:** ≥ 4 , $E_T > 15\text{ GeV}$, $|\eta| < 2$
- **Missing E_T > 20 GeV**
- " E_T^W " > 60 GeV ; $|\eta_W| < 2$

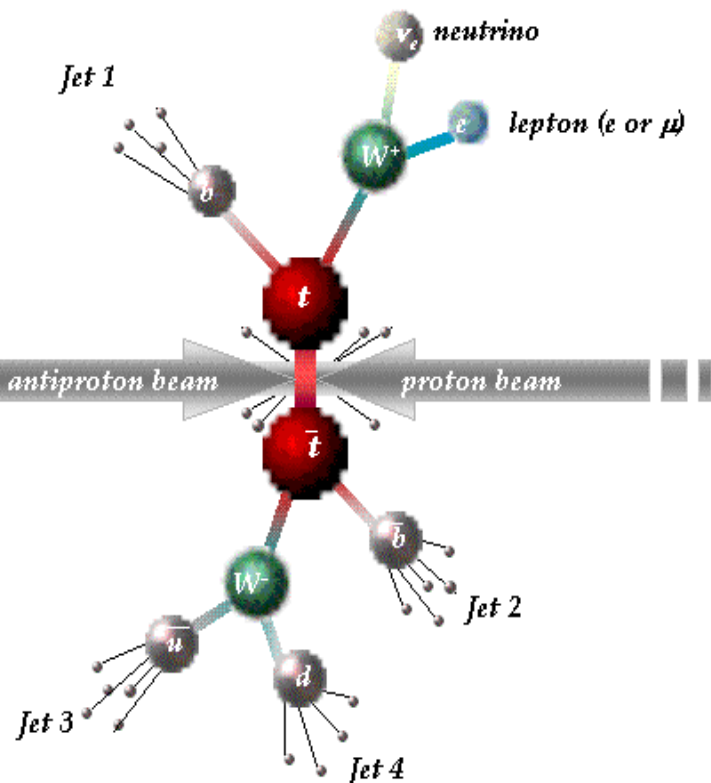
91 events Ref. PRD 58 (1998), 052001

After χ^2 cut 40% signal

Additional cuts for this analysis :

- 4 Jets only (Leading Order Matrix Element)
- **71 events**
- Background probability (to improve purity)

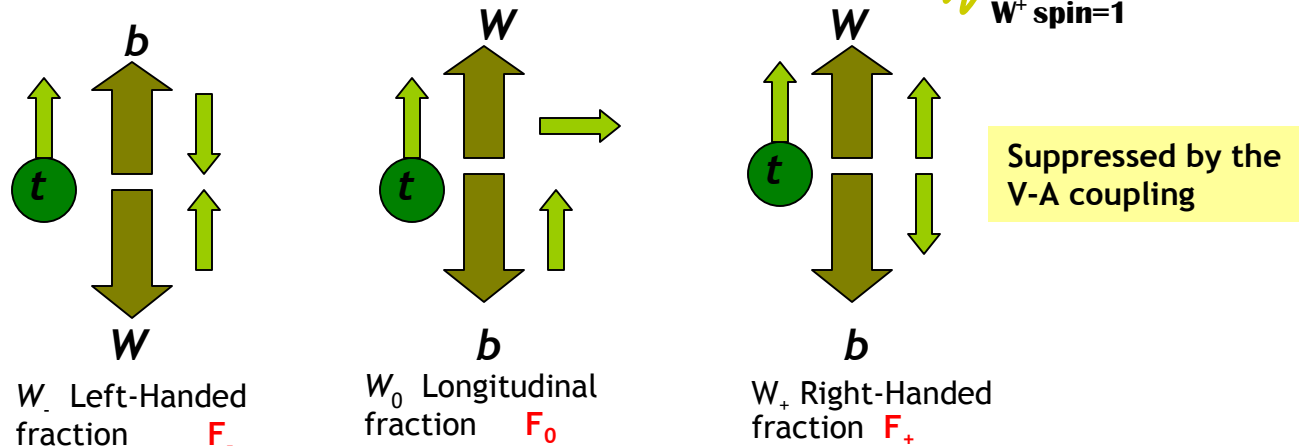
22 events \Rightarrow 12 signal + 10 background



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Helicity of the W in ttbar events

- Top Standard Model weak decay -> V-A coupling as it is for all the other fermions



$$w(\cos \varphi_{l\bar{b}}) = F_- \cdot \frac{3}{8}(1 - \cos \varphi_{l\bar{b}})^2 + F_0 \cdot \frac{3}{8}(1 - \cos^2 \varphi_{l\bar{b}}) + F_+ \cdot \frac{3}{8}(1 + \cos \varphi_{l\bar{b}})^2$$

In SM (with $m_b=0$, $M_{top} = 175$ GeV and $m_W = 80.4$ GeV),

$$F_- = \frac{2 \frac{m_W^2}{M_{top}^2}}{1 + 2 \frac{m_W^2}{M_{top}^2}} \approx 0.30$$

We want to extract

$$F_0 = \frac{1}{1 + 2 \frac{m_W^2}{M_{top}^2}} \approx 0.70$$

$$F_+ = 0$$

The General Method

- We want to find the value of a parameter α

In our case $\alpha = F_0$

- The best estimate of a **parameter** (α) is achieved comparing the events with the probability from the theory with the data. This is done by maximizing a likelihood:

$$L(\alpha) = e^{-N \int \bar{p}(x; \alpha) dx} \prod_{i=1}^N \bar{P}(x_i; \alpha)$$

where x is a set of measured variables

- Probability.** Sum over all the possible parton variables y leading to the observed set of variables x

$d^n \sigma$ is the differential cross section.
ME: F_0 in leptonic and hadronic decays

$W(y, x)$ is the probability that a parton level set of variables y will be measured as a set of variables x

$$\bar{P}(x; \alpha) = \frac{1}{\sigma} \int d^n \sigma(y; \alpha) \underbrace{dq_1 dq_2 f(q_1) f(q_2) W(x, y)}_{f(q) \text{ is the probability distribution than a parton will have a momentum } q}$$

- Detector effects**

$$\bar{P}_{measured}(x; \alpha) = Acc(x) \bar{P}_{production}(x; \alpha)$$

where $Acc(x)$ include all conditions for accepting or rejecting an event

- Background** events with weights c_i $\bar{P}(x; c_1, \dots, c_K, \alpha) = \sum_{i=1}^K c_i \bar{P}_i(x; \alpha)$

Transfer Function $W(x,y)$

- $W(x,y)$ probability of measuring x when y was produced (x jet variables, y parton variables):

Energy of **electrons** is considered well measured

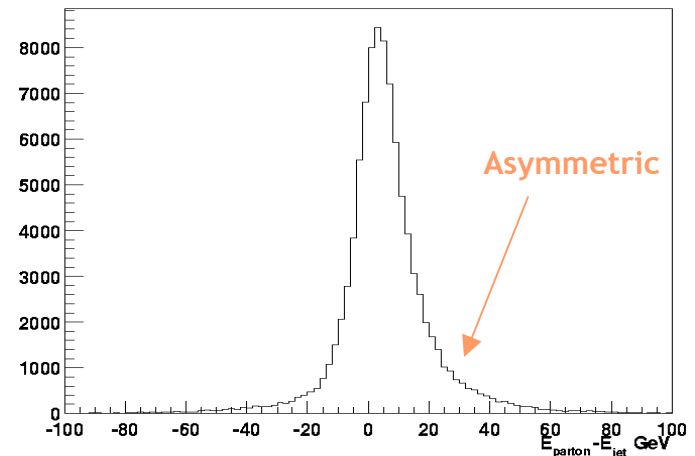
$$W(x,y) = \delta^3(p_e^y - p_e^x) \prod_{j=1}^4 W_{jet}(E_j^y, E_j^x) \prod_{i=1}^4 \delta^2(\Omega_i^y - \Omega_i^x)$$

And due to the excellent granularity of the DØ calorimeter, **angles** are also considered well measured

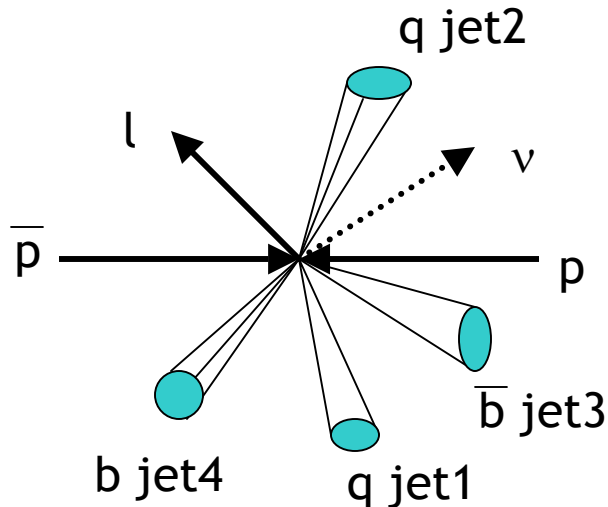
- $W_{jet}(x,y)$ model the smearing in jet energies from effects of **radiation, hadronization, measurement resolution, and jet reconstruction algorithm**

- o Use 2 gaussians, one to account for the peak and the other to fit the asymmetric tails
- o Correcting on average, and considering these distributions to be Gaussian can underestimate the jet energy
- o b and light quarks parameterizations

- Events with muons are integrated over their resolution



Probability for Signal Events



- $2(in) + 18(final) = 20$ degrees of freedom
- $3(e) + 8(\Omega_1 \dots \Omega_4) + 3(P_{in} = P_{final}) + 1(E_{in} = E_{final}) = 15$ constraints
- $20 - 15 = 5$ integrals \Rightarrow we choose M_{top} , m_W and jet energy of one of the jets because $|M|^2$ is almost negligible, except near the four peaks of the Breit-Wigners within $|M|^2$
- All the neutrino all possible solutions are considered
- Sum over 12 combinations of jets

$$P_{t\bar{t}}(x, \alpha) = \frac{1}{12\sigma_{t\bar{t}}} \int d\rho_1 dm_1^2 dM_1^2 dm_2^2 dM_2^2 \sum_{comb, v} |M_{t\bar{t}}(\alpha)|^2 \frac{f(q_1)f(q_2)}{|q_1||q_2|} \phi_6 W_{jet}(x, y)$$

- ρ_1 momentum of one of the jets
- m_1, m_2 top mass in the event
- M_1, M_2 W mass in the event
- $f(q_1), f(q_2)$ parton distribution function (CTEQ4) for incident partons
- q_1, q_2 initial parton momentum
- ϕ_6 six particle phase space
- $W_{jet}(x, y)$ probability of measuring x when y was produced in the collision
- $|M_{t\bar{t}}|^2$ $t\bar{t} \rightarrow \text{lepton_jets}$ matrix element (only qqbar)

Approximations in the probabilities definitions (things to do better with more statistics)

- **Only $t\bar{t}b\bar{a}r$ from $q\bar{q}b\bar{a}r$ production:** it does not include 10% of $t\bar{t}b\bar{a}r$ events that are produced by gluon fusion
- **Only W +jets background:** that is ~85% only of the background
 - ❖ The background probability is defined only in terms of the main background (W +jets, 85%) which proves to be an adequate representation for multijet background
 - ❖ The background probability for each event is calculated using VECBOS subroutines for W +jets
 - ❖ Similar procedure than for $t\bar{t}b\bar{a}r$ events
- **Leading-Order $t\bar{t}b\bar{a}r$ matrix element:** no extra jets, constrains our sample to have only 4 jets

$$P_0(x; c_1, c_2, \alpha) = c_1 P_{t\bar{t}b\bar{a}r}(x; \alpha) + c_2 P_{W+jets}(x)$$

After these approximations, the likelihood function used is

$$-\ln L(\alpha) = -\sum_{i=1}^N \ln [c_1 P_{t\bar{t}b\bar{a}r}(x_i; \alpha) + c_2 P_{W+jets}(x_i)] + N \int A(x) [c_1 P_{t\bar{t}b\bar{a}r}(x; \alpha) + c_2 P_{W+jets}(x)] dx$$

The values of c_1 and c_2 are optimized, and the likelihood is normalized automatically at each value of α

Depends on α

Constant

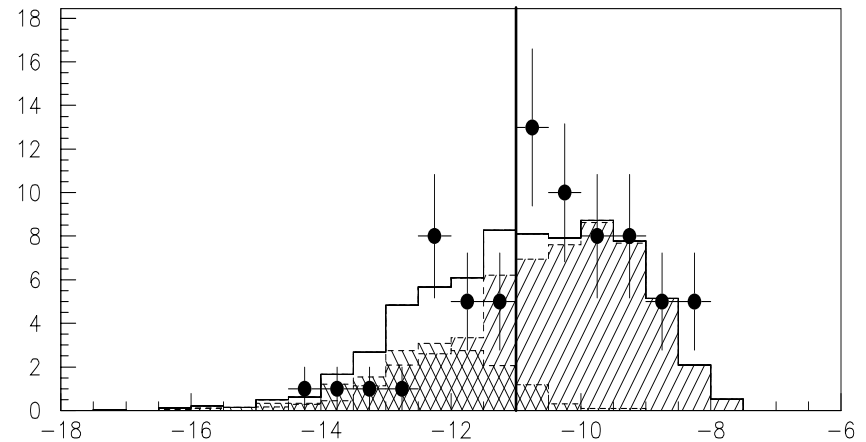
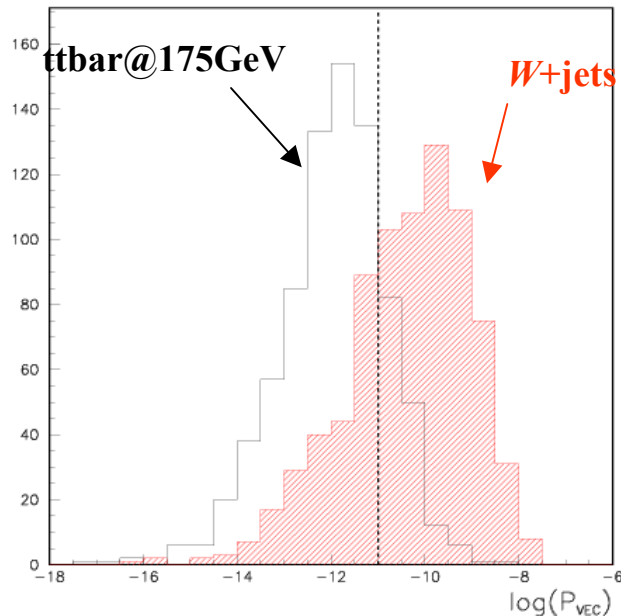
Calculated in two different ways using Monte Carlo method of integration

Extra Selection in P_{bkg}

- In order to increase the purity of signal another selection is applied on P_{bkg} , with efficiencies:

$$\begin{aligned}\epsilon_{t\bar{t}} &= 0.70, \\ \epsilon_{W+jets} &= 0.30\end{aligned}$$

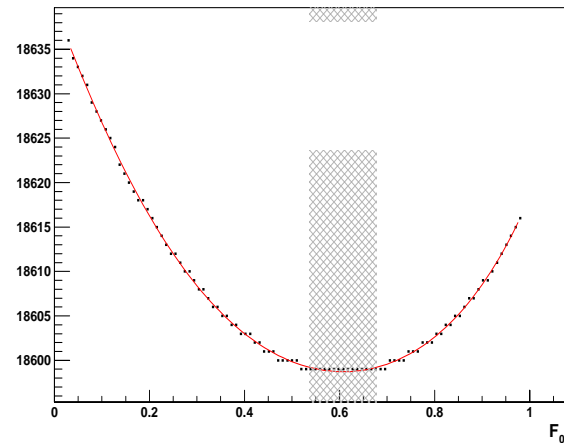
- We select on $P_{bkg} < 10^{-11}$, according to a previous analysis done with this method to measure the top mass



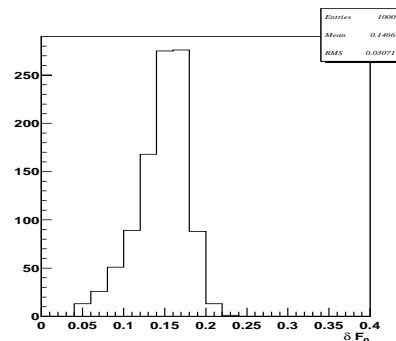
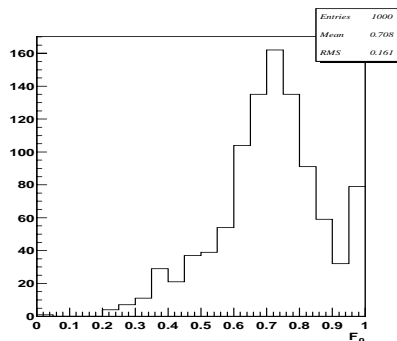
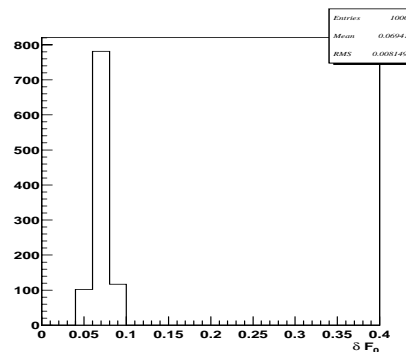
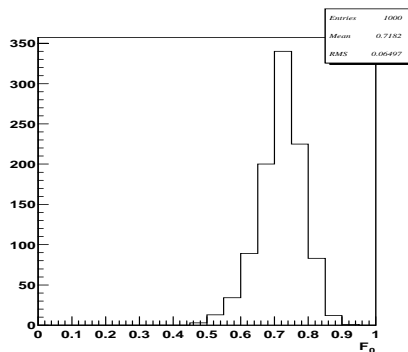
- Comparison of (16 Signal + 55 Background) MC and data sample. Background probability comparison between data (dots) and MC (histogram)

Example using Monte Carlo Events

- Likelihood is calculated in the physical region, 0 to 1
- Since we are dealing with low statistics, we choose to extract the most probable value and its error defined as the $\frac{1}{2}$ 68.27% region around it



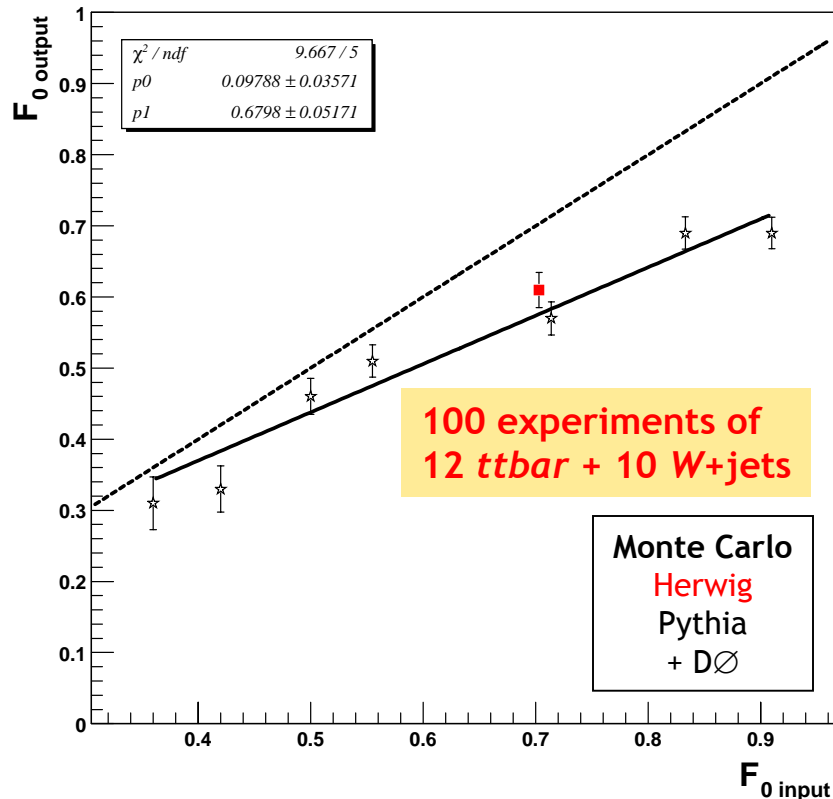
100 events per experiment



20 events per experiment

Linearity of Response for F_0

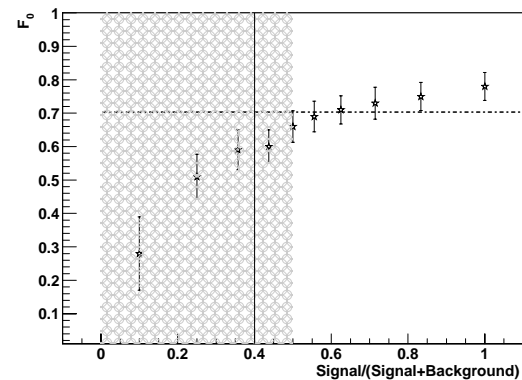
- Using different input values of F_0 , experiments are used to determine output values of F_0
- A response correction needs to be applied to the data



- Output F_0 is biased towards smaller values as more background is introduced

- There is no bias when using parton level $t\bar{t}$ bar and W +jets Monte Carlo events

- Effect may come from radiation

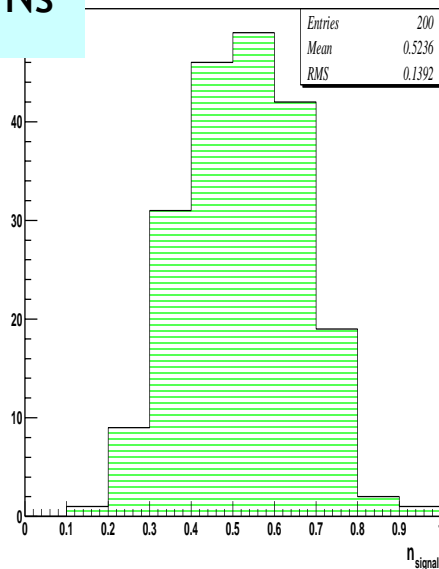


Full MC Simulation

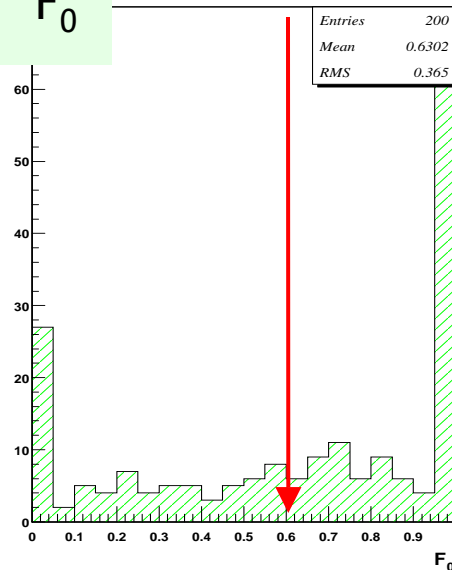
Ensemble Tests (12 Signal + 10 Background)

- 200 experiments of 12 $t\bar{t}$ ($F_0=0.7$) + 10 W +jets events
- Input F_0 is within 68.27% interval of the likelihood in 67% of the experiments
 - reasonable definition for the uncertainty on F_0
- Distributions show most probable F_0 , uncertainty in F_0 , and number of signal events
- **Arrows show Run I data**

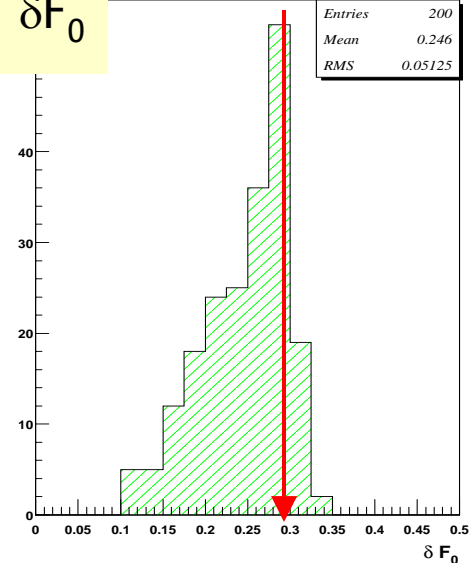
N_s



F_0

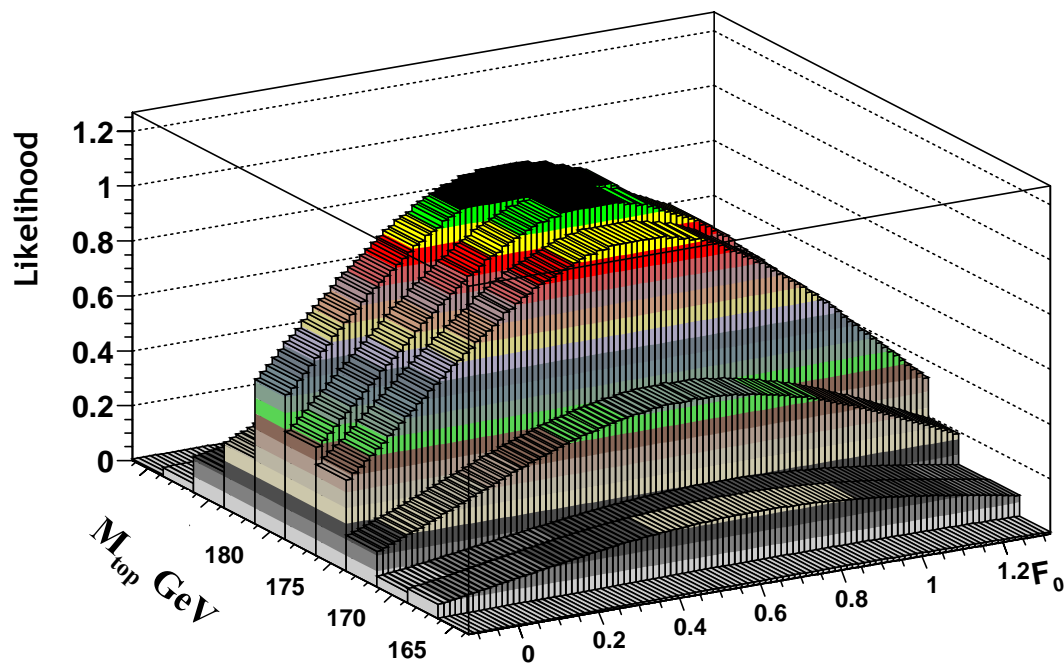


δF_0



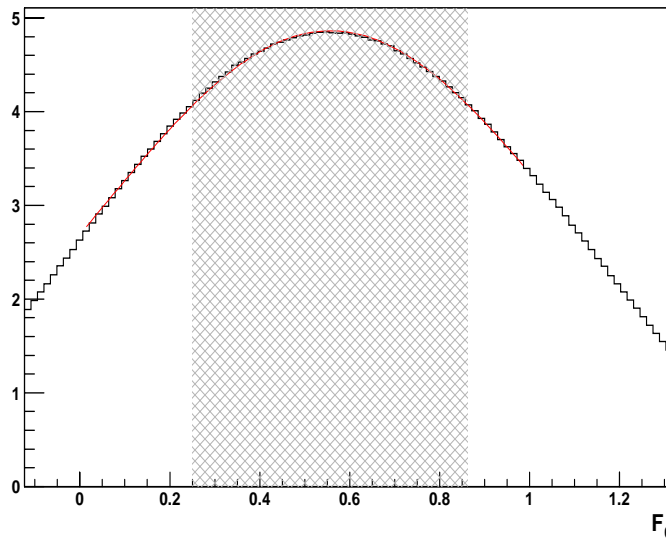
Two-dimensional Probability - M_{top} , F_0

- Assuming $F_0 = 0.7$ (SM), M_{top} is measured to be 180.1 ± 3.6 GeV (shift of 0.5 GeV applied)
- Assuming $M_{top} = 175$ GeV, F_0 is measured to be 0.599 ± 0.302 (linearity response applied)



Preliminary Measurement of F_0 with $D\bar{D}$ Run I Data

- Uncertainty on the top mass translates into a systematic error on the measurement of F_0
- Integrate over M_{top} $L(F_0) = \int L(M_{\text{top}}, F_0) dM_t$
- Most probable value and 68.27% interval using $M_{\text{top}} = 175$ GeV
- 22 events pass our cuts => from fit, **12 signal + 10 background events**



$$F_0 \pm \delta F_0(\text{Stat} + M_{\text{top}}) = 0.558 \pm 0.306$$

From data

Statistics + M_{top}
uncertainty

0.306

Jet Energy Scale

0.014

Parton Distribution
Function

0.007

Acceptance-Linearity
Correction

0.021

From Monte
Carlo

Background

0.010

Signal Model

0.020

Multiple Interactions

0.009

$t\bar{t}$ Spin Correlations

0.008

Conclusions

- ❖ The helicity of the W boson offers a way to learn about the decay coupling of the top quark
- ❖ Using LO approximation and parameterized showering, we calculated the event probabilities, and measured:

$$F_0 \text{ (preliminary)} = 0.56 \pm 0.31$$

First F_0 measurement done at $D\bar{D}$ using **22 events** (~50% signal)

CDF measurement using **108 leptons** (~70% signal) 0.91 ± 0.39

- ❖ This method was first applied to the re-measurement of the **top quark mass**, and now applied to measure **angular distributions**
- ❖ We have a method that allows us to extract F_0 using the **maximal information** in the event:
 - ✓ Correct permutation is always considered (along with the other eleven)
 - ✓ All features of individual events are included, thereby well measured events contribute more information than poorly measured events
 - ✓ This method offers the possibility of increasing the statistics using **both W decay branches**
 - ✓ For higher statistics, one clearly needs to improve the calculation of the probabilities, but this method is a **better way** to do the analysis